## Weyl's Asymptotic Law

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Let  $\Omega$  be a bounded open subset of  $\mathbb{R}^m$  satisfying the segment property, i.e., its topological boundary  $\partial\Omega$  has a finite open covering  $\{O_i\}$  and corresponding nonzero vectors  $\{y_i\}$  such that

$$\{ x + ty_i \mid 0 < t < 1 \} \subset \Omega \,, \qquad \forall x \in \overline{\Omega} \cap O_i \,.$$

We refer to a subset of  $\mathbb{R}^m$  with these properties as a *domain*. Denote by  $C^{\infty}(\overline{\Omega})$  the set of smooth functions on  $\Omega$  all of whose partial derivatives can be extended continuously to  $\overline{\Omega}$ . Consider the subspace

$$D_D = \{ f \in C^{\infty}(\overline{\Omega}) \mid f|_{\partial\Omega} = 0 \}$$

of the Hilbert space  $L^2(\Omega)$ . Let

$$\Delta = -\operatorname{div} \circ \operatorname{grad} \tag{1}$$

be the Laplace operator on  $C^{\infty}(\Omega)$ . If  $\partial\Omega$  is smooth, then, by Green's theorem, the restriction of  $\Delta$  to  $D_D$  is symmetric and non-negative, i.e.,

$$\langle \Delta f, g \rangle = \langle \operatorname{grad} f, \operatorname{grad} g \rangle_{L^2} := \int_{\Omega} (\operatorname{grad} f, \operatorname{grad} g) \, d^m x = \langle f, \Delta g \rangle \,, \quad (2)$$

where  $\langle ., . \rangle$  and (., .) denote the inner product on  $L^2(\Omega)$  and the Euclidean metric on  $\mathbb{R}^m$ , respectively. We denote  $\langle \operatorname{grad} f, \operatorname{grad} g \rangle_{L^2}$  by Q(f, g).

Consider the Sobolev space

$$H^1(\Omega) := W^{1,2}(\Omega) = \left\{ f \in L^2(\Omega) \, | \, \operatorname{Grad} f \in L^2(\Omega) \right\},$$

where Grad f denote the distributional gradient of f. We define the subspace  $H_0^1(\Omega)$  of  $H^1(\Omega)$  to be the closure of the subspace  $C_c^{\infty}(\Omega) \subset H^1(\Omega)$  under the norm  $||.||_{H^1}$ , where  $||f||_{H^1}^2 = ||f||_{L^2}^2 + ||\operatorname{Grad} f||_{L^2}^2$ . The Dirichlet Laplacian for  $\Omega$ , denoted by  $\Delta_D^{\Omega}$ , is defined to be the self-adjoint non-negative Friedrichs

extension of  $\Delta|_{D_D}$ . The domain of  $\Delta_D^{\Omega}$  is the subspace  $H_0^1(\Omega)$  of  $L^2(\Omega)$ . We have

$$\langle \Delta_D^{\Omega} f, f \rangle = \overline{Q}(f, f), \qquad \forall f \in H_0^1(\Omega),$$
(3)

where  $\overline{Q}$  denote the closure of Q.

Consider the subspace  $D_N = \{f \in C^{\infty}(\overline{\Omega}) \mid (vf)|_{\partial\Omega} = 0\}$  of  $L^2(\Omega)$ , where v denote the outward unit normal vector field on  $\partial\Omega$ . If  $\partial\Omega$  is smooth, then, by Green's theorem, the restriction of  $\Delta$  to  $D_N$  is symmetric and non-negative. The Neumann Laplacian for  $\Omega$ , denoted by  $\Delta_N^{\Omega}$ , is defined to be the self-adjoint non-negative Friedrichs extension of  $\Delta|_{D_N}$ . The domain of  $\Delta_N^{\Omega}$  is the subspace  $H^1(\Omega)$  of  $L^2(\Omega)$ . For each  $f \in H^1(\Omega)$  we have  $\langle \Delta_N^{\Omega} f, f \rangle = \overline{Q}(f, f)$ .

The self-adjoint non-negative operators  $\Delta_D^{\Omega}$  and  $\Delta_N^{\Omega}$  have compact resolvent [3]. Thus, their (point) spectrum consist of a discrete increasing sequence of non-negative real numbers with no limit point in  $\mathbb{R}$ , and each eigenvalue has finite multiplicity. In addition, the set of corresponding eigenfunctions is an orthonormal basis for  $L^2(\Omega)$ .

Henceforth, we denote by  $\phi_k^{\Omega}$  (respectively,  $\psi_k^{\Omega}$ ) the normalized eigenfunction of  $\Delta_D^{\Omega}$  (respectively,  $\Delta_N^{\Omega}$ ) with eigenvalue  $\lambda_k^{\Omega}$  (respectively,  $\mu_k^{\Omega}$ ), where k ranges over positive integers, i.e.,  $k = 1, 2, \cdots$ . We arrange the eigenvalues in an increasing sequence

$$0 \le \lambda_1^{\Omega} \le \lambda_2^{\Omega} \le \cdots$$
$$0 \le \mu_1^{\Omega} \le \mu_2^{\Omega} \le \cdots$$

with each eigenvalue repeated according to its multiplicity. For each  $\nu \in \mathbb{R}_{>0}$ , let  $N(\Delta_D^{\Omega}, \nu)$  be the number of eigenvalues  $\lambda_k^{\Omega} \leq \nu$  counted with multiplicity. Indeed,  $N(\Delta_D^{\Omega}, \nu)$  is the sum of the dimensions of the eigenspaces of  $\Delta_D^{\Omega}$ corresponding to the eigenvalues  $\lambda_k^{\Omega} \leq \nu$ . The quantity  $N(\Delta_N^{\Omega}, \nu)$  is defined similarly. In the following, we want to investigate the asymptotic behaviour of  $N(\Delta_D^{\Omega}, \nu)$  as  $\nu \to \infty$ .

As the first step, we need to study the asymptotic behaviour of  $N(\Delta_D^{\Omega}, \nu)$ and  $N(\Delta_N^{\Omega}, \nu)$  in the special case where,  $\Omega$  is the open *m*-dimensional rectangle  $\Omega = (0, l_1) \times (0, l_2) \times \cdots \times (0, l_m) \subset \mathbb{R}^m$ . In this case, the eigenvalues of  $\Delta_D^{\Omega}$  and  $\Delta_N^{\Omega}$  are explicitly given by (see [3], [1]):

$$\lambda^{\Omega} = \sum_{i=1}^{m} k_i^2 \left(\frac{\pi}{l_i}\right)^2 \tag{4}$$

and

$$\mu^{\Omega} = \sum_{i=1}^{m} (k_i - 1)^2 \left(\frac{\pi}{l_i}\right)^2,$$
(5)

respectively, where each  $k_i$  ranges over the positive integers.

**Proposition 1** Let  $\Omega$  be the open *m*-dimensional rectangle as above. We have

$$\lim_{\nu \to \infty} \frac{N(\Delta_D^{\Omega}, \nu)}{\nu^{m/2}} = \lim_{\nu \to \infty} \frac{N(\Delta_N^{\Omega}, \nu)}{\nu^{m/2}} = \frac{\omega_m \operatorname{Vol}(\Omega)}{(2\pi)^m},$$
(6)

where  $\omega_m$  is the volume of the unit ball in  $\mathbb{R}^m$ , and  $Vol(\Omega)$  denotes the volume of  $\Omega$ .

**Proof** Let  $\Gamma$  be the lattice in  $\mathbb{R}^m$  generated by the vectors  $\{e_1/l_1, \dots, e_m/l_m\}$ , where  $\{e_i\}$  are the standard basis of  $\mathbb{R}^m$ . Let  $\mathbb{B}^m(r) \subset \mathbb{R}^m$  be the closed ball with radius r and centered at the origin. Considering (4), the quantity  $N(\Delta_D^{\Omega}, \nu)$  equals the number of the points of  $\Gamma$  in an "octant" of  $\mathbb{B}^m(\sqrt{\nu}/\pi)$ .

Denote by  $\mathcal{N}(r)$  the number of the points of  $\Gamma$  in  $\mathbb{B}^m(r)$ . Let  $\Sigma$  be the fundamental domain for the action of  $\Gamma$  on  $\mathbb{R}^m$ . Denote by  $\mathcal{P}(r)$  the number of copies of  $\Sigma$  contained in  $\mathbb{B}^m(r)$ . Let  $d = \sqrt{\sum_{i=1}^m 1/l_i^2}$  be the diameter of  $\Sigma$ . We have

$$\mathcal{P}(r) \le \mathcal{N}(r) \le \mathcal{P}(r+d)$$
. (7)

Thus, using

$$\omega_m (r-d)^m \le \mathcal{P}(r) \operatorname{Vol}(\Sigma) \le \omega_m r^m \,, \tag{8}$$

we get

$$\frac{\omega_m (r-d)^m}{\operatorname{Vol}(\Sigma)} \le \mathcal{N}(r) \le \frac{\omega_m (r+d)^m}{\operatorname{Vol}(\Sigma)},\tag{9}$$

which implies

$$\lim_{r \to \infty} \frac{\mathcal{N}(r)}{r^m} = \frac{\omega_m}{\operatorname{Vol}(\Sigma)} = \omega_m \operatorname{Vol}(\Omega) \,. \tag{10}$$

So we have  $N(\Delta_D^{\Omega}, \nu) = \mathcal{N}(\sqrt{\nu}/\pi)/2^m + \epsilon$ , where the "error" term  $\epsilon$  corresponds to the points of the (m-1)-dimensional lattices resulted from intersection of  $\Gamma$  with each coordinate hyperplane in  $\mathbb{R}^m$ . Since  $\epsilon$  is of order  $\nu^{(m-1)/2}$ , using (10), we get (6) for  $N(\Delta_D^{\Omega}, \nu)$ . A similar argument yields to (6) for  $N(\Delta_N^{\Omega}, \nu)$ .

The next step is to show the *domain monotonicity* of the eigenvalues of  $\Delta_D$  and  $\Delta_N$ , and also, find a relation between the corresponding eigenvalues of  $\Delta_D^{\Omega}$  and  $\Delta_N^{\Omega}$  on the same domain  $\Omega$ .

We recall the max-min principle for the eigenvalues of self-adjoint positive operators. Consider a domain  $\Omega$  in  $\mathbb{R}^m$ . Let  $E(f) := \overline{Q}(f, f)/||f||_{L^2}^2$  be the Rayleigh quotient. Denote by  $\langle g_1, \dots, g_n \rangle$  the span of the the functions  $g_i \in L^2(\Omega), i = 1, \dots, n$ . According to the max-min principle, we have

$$\lambda_k^{\Omega} \le E(f) \,, \qquad \forall f \in \langle \phi_1^{\Omega}, \cdots, \phi_{k-1}^{\Omega} \rangle^{\perp} \cap H_0^1(\Omega) \,, \tag{11}$$

$$\mu_k^{\Omega} \le E(f), \quad \forall f \in \langle \psi_1^{\Omega}, \cdots, \psi_{k-1}^{\Omega} \rangle^{\perp} \cap H^1(\Omega).$$
(12)

**Lemma 2** Consider the domains  $\tilde{\Omega}$  and  $\Omega$ , where  $\tilde{\Omega} \subset \Omega \subset \mathbb{R}^m$ . We have

$$\lambda_k^{\Omega} \le \lambda_k^{\tilde{\Omega}} \tag{13}$$

for each  $k = 1, 2, \cdots$ .

**Proof** Note that each  $f \in L^2(\tilde{\Omega})$  can be viewed as as an element of  $L^2(\Omega)$  by setting it equal to zero in  $\Omega \setminus \tilde{\Omega}$ . Consider the linear map

$$T: \langle \phi_1^{\tilde{\Omega}}, \cdots, \phi_k^{\tilde{\Omega}} \rangle \to \langle \phi_1^{\Omega}, \cdots, \phi_{k-1}^{\Omega} \rangle^*$$
$$\phi_i^{\tilde{\Omega}} \mapsto \langle ., \phi_i^{\tilde{\Omega}} \rangle, \qquad i = 1, \cdots, k$$

Since  $\operatorname{Ker}(T) \neq \emptyset$ , there exist a function  $f_0 = \sum_{i=1}^k a_i \phi_i^{\tilde{\Omega}}$ ,  $a_i \in \mathbb{R}$  such that  $f_0 \in \langle \phi_1^{\Omega}, \cdots, \phi_{k-1}^{\Omega} \rangle^{\perp} \cap H_0^1(\Omega)$ . Considering (11), we have

$$\lambda_{k}^{\Omega} \leq E(f_{0}) = \frac{\langle \Delta_{D}^{\tilde{\Omega}} f_{0}, f_{0} \rangle}{||f_{0}||_{L^{2}}^{2}} = \frac{\sum_{i=1}^{k} \lambda_{i}^{\tilde{\Omega}} a_{i}^{2}}{\sum_{i=1}^{k} a_{i}^{2}} \leq \lambda_{k}^{\tilde{\Omega}}.$$
 (14)

**Lemma 3** Let  $\Omega$  be a domain in  $\mathbb{R}^m$ . We have

$$\mu_k^{\Omega} \le \lambda_k^{\Omega} \tag{15}$$

for each  $k = 1, 2, \cdots$ .

**Proof** Since  $H_0^1(\Omega) \subset H^1(\Omega)$ , there exist a function  $f_0 = \sum_{i=1}^k a_i \phi_i^{\Omega}$ ,  $a_i \in \mathbb{R}$  such that  $f_0 \in \langle \psi_1^{\Omega}, \cdots, \psi_{k-1}^{\Omega} \rangle^{\perp} \cap H^1(\Omega)$ . Considering (12), we have

$$\mu_k^{\Omega} \le E(f_0) = \frac{\langle \Delta_D^{\Omega} f_0, f_0 \rangle}{||f_0||_{L^2}^2} = \frac{\sum_{i=1}^k \lambda_i^{\Omega} a_i^2}{\sum_{i=1}^k a_i^2} \le \lambda_k^{\Omega} .$$
(16)

Consider the disjoint domains  $\Omega_1, \Omega_2$  in the domain  $\Omega \subset \mathbb{R}^m$  such that  $\Omega = int(\Omega_1 \cup \Omega_2)$ , and  $\Omega \setminus \Omega_1 \cup \Omega_2$  has Lebesgue measure zero. (Intuitively, one gets  $\Omega_1 \cup \Omega_2$  by eliminating a "surface"  $(\Omega \setminus \Omega_1 \cup \Omega_2) \subset int\Omega$  from  $\Omega$ ). Note that, since  $H^1(\Omega_1 \cup \Omega_2) = H^1(\Omega_1) \oplus H^1(\Omega_2)$ , we have

$$\Delta_D^{\Omega_1 \cup \Omega_2} = \Delta_D^{\Omega_1} \oplus \Delta_D^{\Omega_2}$$
$$\Delta_N^{\Omega_1 \cup \Omega_2} = \Delta_N^{\Omega_1} \oplus \Delta_N^{\Omega_2}.$$
 (17)

**Lemma 4** Let  $\Omega_1, \Omega_2 \subset \Omega$  be as above. We have

$$\mu_k^{\Omega_1 \cup \Omega_2} \le \mu_k^{\Omega} \,, \tag{18}$$

$$\lambda_k^{\Omega} \le \lambda_k^{\Omega_1 \cup \Omega_2} \tag{19}$$

for each  $k = 1, 2, \cdots$ .

**Proof** If  $f \in H^1(\Omega)$ , its restriction to  $\Omega_1 \cup \Omega_2$  is in  $H^1(\Omega_1) \oplus H^1(\Omega_2)$ . So there exist a function  $f_0 = \sum_{i=1}^k a_i \psi_i^{\Omega}$ ,  $a_i \in \mathbb{R}$  such that  $f_0 \in \langle \psi_1^{\Omega_1 \cup \Omega_2}, \cdots, \psi_{k-1}^{\Omega_1 \cup \Omega_2} \rangle^{\perp} \cap H^1(\Omega_1 \cup \Omega_2)$ . In addition, since  $\Omega \setminus \Omega_1 \cup \Omega_2$  has

Lebesgue measure zero, we have

$$\int_{\Omega_1 \cup \Omega_2} \left| \operatorname{Grad} f_0 \right|^2 d^m x = \int_{\Omega} \left| \operatorname{Grad} f_0 \right|^2 d^m x \,. \tag{20}$$

Thus, considering (12), we get

$$\mu_k^{\Omega_1 \cup \Omega_2} ||f_0||_{L^2}^2 \leq \int_{\Omega_1 \cup \Omega_2} |\operatorname{Grad} f_0|^2 d^m x = \int_{\Omega} |\operatorname{Grad} f_0|^2 d^m x$$
$$= \langle \Delta_N^{\Omega} f_0, f_0 \rangle$$
$$= \sum_{i=1}^k \mu_i^{\Omega} a_i^2 \leq \mu_k^{\Omega} ||f_0||_{L^2}^2.$$
(21)

The inequality (19) is a special case of (13).

**Definition** A standard  $2^{-n}$  cube in  $\mathbb{R}^m$  is a cube of the form

$$\left[\frac{b_1}{2^n},\frac{b_1+1}{2^n}\right)\times\cdots\times\left[\frac{b_m}{2^n},\frac{b_m+1}{2^n}\right)$$

with  $b_1, \dots, b_m$  integers. Given a Lebesgue measurable set  $\Omega \subset \mathbb{R}^m$ , denote by  $\Omega_n^-$  (respectively,  $\Omega_n^+$ ) those standard  $2^{-n}$  cubes contained in  $\Omega$  (respectively, those standard  $2^{-n}$  cubes that intersect  $\Omega$ ). Let  $W_n^{\pm}(\Omega) = \operatorname{Vol}(\Omega_n^{\pm})$ . The set  $\Omega$  is called a *contented (Jordan measurable)* set if we have

$$\lim_{n \to \infty} W_n^-(\Omega) = \lim_{n \to \infty} W_n^+(\Omega) = \operatorname{Vol}(\Omega) \,. \tag{22}$$

**Theorem 5** Let  $\Omega$  be a contented domain in  $\mathbb{R}^m$ . We have

$$\lim_{\nu \to \infty} \frac{N(\Delta_D^{\Omega}, \nu)}{\nu^{m/2}} = \frac{\omega_m \operatorname{Vol}(\Omega)}{(2\pi)^m} \,.$$
(23)

**Proof** Let  $\{C_{n,j}^{\pm}\}$  be the interiors of the cubes in the definition of  $\Omega_n^{\pm}$  so that  $\bar{\Omega}_n^{\pm} = \bigcup_j \bar{C}_{n,j}^{\pm}$ . By (13) and (19) we have

$$\lambda_k^{\Omega} \le \lambda_k^{\Omega_n^-} \le \lambda_k^{\cup_j C_{n,j}^-} \,, \tag{24}$$

which implies

$$N(\Delta_D^{\Omega},\nu) \ge N(\Delta_D^{\Omega_n^-},\nu) \ge N(\oplus_j \Delta_D^{C_{n,j}^-},\nu) = \sum_j N(\Delta_D^{C_{n,j}^-},\nu).$$
(25)

Thus, by (6), we have

$$\liminf_{\nu \to \infty} \frac{N(\Delta_D^{\Omega}, \nu)}{\nu^{m/2}} \ge \sum_j \lim_{\nu \to \infty} \frac{N(\Delta_D^{C_{n,j}}, \nu)}{\nu^{m/2}} = \sum_j \frac{\omega_m \operatorname{Vol}\left(C_{n,j}^-\right)}{(2\pi)^m} = \frac{\omega_m W_n^-(\Omega)}{(2\pi)^m}.$$
 (26)

On the other hand, by (15) and (18), we have

$$\lambda_k^{\Omega} \ge \lambda_k^{\Omega_n^+} \ge \mu_k^{\Omega_n^+} \ge \mu_k^{\cup_j C_{n,j}^+}, \qquad (27)$$

which implies

$$N(\Delta_D^{\Omega}, \nu) \le \sum_j N(\Delta_N^{C_{n,j}^+}, \nu) \,. \tag{28}$$

Therefore, by (6), we have

$$\limsup_{\nu \to \infty} \frac{N(\Delta_D^{\Omega}, \nu)}{\nu^{m/2}} \le \sum_j \lim_{\nu \to \infty} \frac{N(\Delta_N^{C_{n,j}^+}, \nu)}{\nu^{m/2}} = \frac{\omega_m W_n^+(\Omega)}{(2\pi)^m} \,. \tag{29}$$

Since  $\Omega$  is contented, if we let  $n \to \infty$ , then, using (26) and (29), we get (23).

## References

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